# Black holes and mathematical sandpaper



Jens Boos boos@ualberta.ca University of Alberta

Joint work with Valeri Frolov and Andrei Zelnikov.

Tenth Annual Symposium for Graduate Physics Research, Thursday, September 26, 2019, 09:45am

## Rotating black holes as a testing ground for new physics



singularities

A singularity is the end of the line where space and time end.

Mathematically: not part of manifold.

Practically: the spacetime curvature grows to infinity, for example:

 $R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}\sim \frac{m^2}{r^6}$ 

2/7

A singularity is the end of the line where space and time end.

Mathematically: not part of manifold.

Practically: the spacetime curvature grows to infinity, for example:

 $R_{\mu\nu
ho\sigma}R^{\mu
u
ho\sigma}\sim rac{m^2}{r^6}$ 

ALLA D

Parter

8 C

2/7

A singularity is the end of the line where space and time end.

Mathematically: not part of manifold.

Practically: the spacetime curvature grows to infinity, for example:

 $R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}\sim \frac{m^2}{r^6}$ 

Parton

5 -

A singularity is the end of the line where space and time end.

Mathematically: not part of manifold.

 $\sigma$ 

m

Entre

Practically: the spacetime curvatur

grows to infir

AAAAA

 $\mu\nu\rho$ 

... but how?

----

#### A simple example: the gravitational potential of a point mass

The gravitational potential of a point-particle is singular at the origin:

$$\nabla^2 \phi = 4\pi Gm\delta(\mathbf{r}) \qquad \longrightarrow \qquad \phi(r) = \frac{-Gm}{r}$$

#### A simple example: the gravitational potential of a point mass

The gravitational potential of a point-particle is singular at the origin:

$$\nabla^2 \phi = 4\pi Gm\delta(\mathbf{r}) \qquad \longrightarrow \qquad \phi(r) = \frac{-Gm}{r}$$

Modify the Poisson equation by an operator  $\hat{F}$  that contains **infinitely many derivatives**:

$$\hat{F} \nabla^2 \phi = 4\pi Gm\delta(\mathbf{r}), \qquad \hat{F} = \exp(\ell^2 \nabla^2) = \sum_{n=0}^{\infty} \frac{1}{n!} \left(\ell^2 \nabla^2\right)^n$$

#### A simple example: the gravitational potential of a point mass

The gravitational potential of a point-particle is singular at the origin:

$$\nabla^2 \phi = 4\pi Gm\delta(\mathbf{r}) \qquad \longrightarrow \qquad \phi(r) = \frac{-Gm}{r}$$

Modify the Poisson equation by an operator  $\hat{F}$  that contains **infinitely many derivatives**:

$$\hat{F} \nabla^2 \phi = 4\pi Gm\delta(\mathbf{r}), \qquad \hat{F} = \exp(\ell^2 \nabla^2) = \sum_{n=0}^{\infty} \frac{1}{n!} \left(\ell^2 \nabla^2\right)^n$$

Okay, but why does that solve our problem? Let's do a simple "back of the envelope" calculation!

Let's consider a simple one-dimensional example:

 $\hat{F} \nabla^2 \phi = 4\pi Gm\delta(r)$ 

Let's consider a simple one-dimensional example:

$$\nabla^2 \phi = 4\pi G m \hat{F}^{-1} \delta(r)$$

Let's consider a simple one-dimensional example:

$$\nabla^2 \phi = 4\pi G m \hat{F}^{-1} \delta(r) = 4\pi G m \,\delta_{\text{eff}}(r)$$

Let's consider a simple one-dimensional example:

$$\nabla^2 \phi = 4\pi G m \hat{F}^{-1} \delta(r) = 4\pi G m \,\delta_{\text{eff}}(r)$$

 $\delta_{\rm eff}(r) = \hat{F}^{-1}\delta(r)$ 

Let's consider a simple one-dimensional example:

$$\nabla^2 \phi = 4\pi G m \hat{F}^{-1} \delta(r) = 4\pi G m \,\delta_{\text{eff}}(r)$$

$$\delta_{\text{eff}}(r) = \hat{F}^{-1}\delta(r) = \exp(\ell^2 \partial_r^2) \int_{-\infty}^{\infty} \frac{\mathrm{d}k}{2\pi} e^{ikr} = \int_{-\infty}^{\infty} \frac{\mathrm{d}k}{2\pi} e^{-k^2 \ell^2 + ikr} = \frac{e^{-r^2/(4\ell^2)}}{\sqrt{4\pi}\ell} \,.$$

Let's consider a simple one-dimensional example:

$$\nabla^2 \phi = 4\pi G m \hat{F}^{-1} \delta(r) = 4\pi G m \,\delta_{\text{eff}}(r)$$
  
$$\tilde{b}_{\text{eff}}(r) = \hat{F}^{-1} \delta(r)$$
  
$$= \exp(\ell^2 \partial_r^2) \int_{-\infty}^{\infty} \frac{\mathrm{d}k}{2\pi} e^{ikr} = \int_{-\infty}^{\infty} \frac{\mathrm{d}k}{2\pi} e^{-k^2 \ell^2 + ikr} = \frac{e^{-r^2/(4\ell^2)}}{\sqrt{4\pi}\ell}$$

What have we done? The delta-shaped source has been smoothed! The result is just a Gaussian bell curve.

Let's consider a simple one-dimensional example:

 $\delta$ 

$$\nabla^2 \phi = 4\pi G m \hat{F}^{-1} \delta(r) = 4\pi G m \,\delta_{\text{eff}}(r)$$

$$eff(r) = \hat{F}^{-1}\delta(r)$$

$$= \exp(\ell^2 \partial_r^2) \int_{-\infty}^{\infty} \frac{\mathrm{d}k}{2\pi} e^{ikr} = \int_{-\infty}^{\infty} \frac{\mathrm{d}k}{2\pi} e^{-k^2 \ell^2 + ikr} = \frac{e^{-r^2/(4\ell^2)}}{\sqrt{4\pi\ell}}$$

What have we done? The delta-shaped source has been smoothed! The result is just a Gaussian bell curve.

#### The delta-shaped source has been smoothed!



Belt grinder, Matthias Wandel (2018).

Sugar 2

400

#### Application to black holes?

Infinite-derivative operators ("mathematical sandpaper") can be used to smooth sharp, local quantities over a characteristic length scale  $\ell$ . This works well for linear equations. Not linear:

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^3} T_{\mu\nu}$$

## Application to black holes?

Infinite-derivative operators ("mathematical sandpaper") can be used to smooth sharp, local quantities over a characteristic length scale  $\ell$ . This works well for linear equations. Not linear:



## Application to black holes?

Infinite-derivative operators ("mathematical sandpaper") can be used to smooth sharp, local quantities over a characteristic length scale  $\ell$ . This works well for linear equations. Not linear:



They are ten non-linear partial differential equations for the ten functions  $g_{\mu\nu}(x)$ , if one provides the energy-momentum (energy density, pressure, shear,...)  $T_{\mu\nu}(x)$ .

#### Thank you for your attention!

#### Abstract

In my ten minutes I would like to present the idea of smoothing black hole singularities by using infinite-derivative operators. I will show in a simple calculation how these operators indeed smooth delta-shaped potentials and act as "mathematical sandpaper." Applying this to black holes is difficult, but I will try to outline the main conceptual steps on the journey towards singularity-free black holes.